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A PROBLEM IN THE HEATING OF SMALL AREAS ON A LARGE SOLID

by

P. H. Thomas

Summary

The transient temperature at the centre of a heated circular area on a semi-infinite solid with cooling to the atmosphere has been computed.

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Fire Research Station,  
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Introduction.

In F.R. Note 163/1955 (1) Blok's (2) analysis for the transient heating of a circular area without cooling was given. Steady state temperatures with cooling were also derived. The present note gives the temperature at the centre of circular area at various times with different degrees of cooling. The work is relevant to the use of small irradiated areas to represent large areas and allows some estimate to be made of the time for which the small area behaves as a large one.

Theory

- Let  $\theta$  be the temperature
- I the incident flux
- H the cooling coefficient
- K the thermal conductivity
- h the quotient  $H/K$
- $\rho$  the density
- c the specific heat
- k the thermal diffusivity  $K/\rho c$
- $r, z$  cylindrical coordinates
- x y z cartesian coordination
- t the time
- and R the radius of heated area

We make use of Green's function for a point some at  $x' y' z'$  in a semi-infinite solid ( $x > 0$ ) with radiation at the boundary  $x = 0$  to a medium at zero temperature. This function is given by Carslaw and Jaeger in "The Conduction of heat in solids" (3). If u is the temperature due to an instantaneous point source of unit strength at  $x', y', z'$ , the temperature at  $x, y, z$  after a time 't' is

$$u = \frac{1}{8(\pi R t)^{3/2}} \left\{ e^{-\frac{(x+x')^2}{4Rt}} + e^{-\frac{(x-x')^2}{4Rt}} \right\} e^{-\frac{((y-y')^2 + (z-z')^2)}{4Rt}}$$

$$- \frac{h}{4\pi R t} e^{-h^2 R t} \left( \frac{x+y'}{2\sqrt{Rt}} + h\sqrt{Rt} \right) x \cdot x \exp \left[ h(x+x') + h^2 R t - \frac{(y-y')^2 + (z-z')^2}{4Rt} \right] \dots (1)$$

For a continuous source of strength  $\frac{I}{\rho c}$  (over  $\sqrt{x^2 + y^2}$ ) at  $x' = 0$ , and  $\sqrt{x^2 + y^2}$  equal to r the temperature at  $x = y = z = 0$  at time 't' is, from (1).

$$\theta = \frac{I}{\rho c} \int_0^t d\lambda \int_0^R r dr e^{-\frac{r^2}{4R\lambda}} \left\{ \frac{1}{(4(\pi R \lambda))^{3/2}} - \frac{h}{4\pi R \lambda} e^{-h^2 R \lambda} e^{-h\sqrt{R\lambda}} \right\} \dots (2)$$

$$= \frac{I}{\rho c} \int_0^t d\lambda (1 - e^{-\frac{R^2}{4R\lambda}}) \left( \frac{1}{(4(\pi R \lambda))^{3/2}} - h e^{-h^2 R \lambda} e^{-h\sqrt{R\lambda}} \right) \dots (3)$$

The solution to the integral of the first term in the first bracket in equation (3) is the result for a heated disk with no surface cooling and is equal to Blok's (2) result

$$\theta_{H=0} = \frac{IR}{K} \left\{ \frac{2}{\sqrt{\pi}} \left( \frac{Rt}{R^2} \right)^{1/2} \left( 1 - e^{-\frac{R^2}{4Rt}} \right) + \operatorname{erfc} \frac{R}{2\sqrt{Rt}} \right\} \dots\dots (4)$$

Hence

$$\theta_{H=0} - \theta = \frac{IR}{\rho c} \int_0^t d\lambda \cdot e^{-h^2 R \lambda} \operatorname{erfc} R \sqrt{R \lambda} \left( 1 - e^{-\frac{R^2}{4R \lambda}} \right) \dots\dots (5)$$

Let  $\lambda = \frac{x^2}{h^2 R}$

Then

$$\theta_{H=0} - \theta = \frac{IR}{K} \cdot \frac{2}{hR} \int_0^{R\sqrt{Rt}} \left( 1 - e^{-\frac{R^2 R^2}{4x^2}} \right) x e^{-\frac{x^2}{R}} \operatorname{erfc} x dx \dots\dots (6)$$

This has been numerically evaluated and in Fig. (1) are shown the resulting values of  $\theta$  for various value of  $hR$  and  $Rt/R^2$ . For the steady state the value of  $\theta$  has been obtained analytically (1) as

$$\theta = \frac{IR}{K} \left[ 1 + \frac{1}{hR} - \frac{\pi}{2} \left( H_1(hR) - Y_1(hR) \right) \right] \dots\dots (7)$$

Where  $Y_1$  is Weber's Bessel function of the second kind of the first order and  $H_1$  is Struve's function of the first order.

It is possible to find the values of  $Rt/R^2$  at which the temperature at the centre of the circle is, say, 10% and 20% less than it would be for a plane source. These values of  $Rt/R^2$  are shown in Fig. (2) as a function of  $hR$ , while in Fig. (3)  $h^2 R t$  is shown as a function of  $hR$  for 10% & 20% error. For any  $h, R,$  and  $k$  it is thus possible to find the maximum duration of heating for which a small source is sensibly similar to a large source.

Acknowledgment

The author is indebted to Mr. Philips who performed the bulk of the computations.

References

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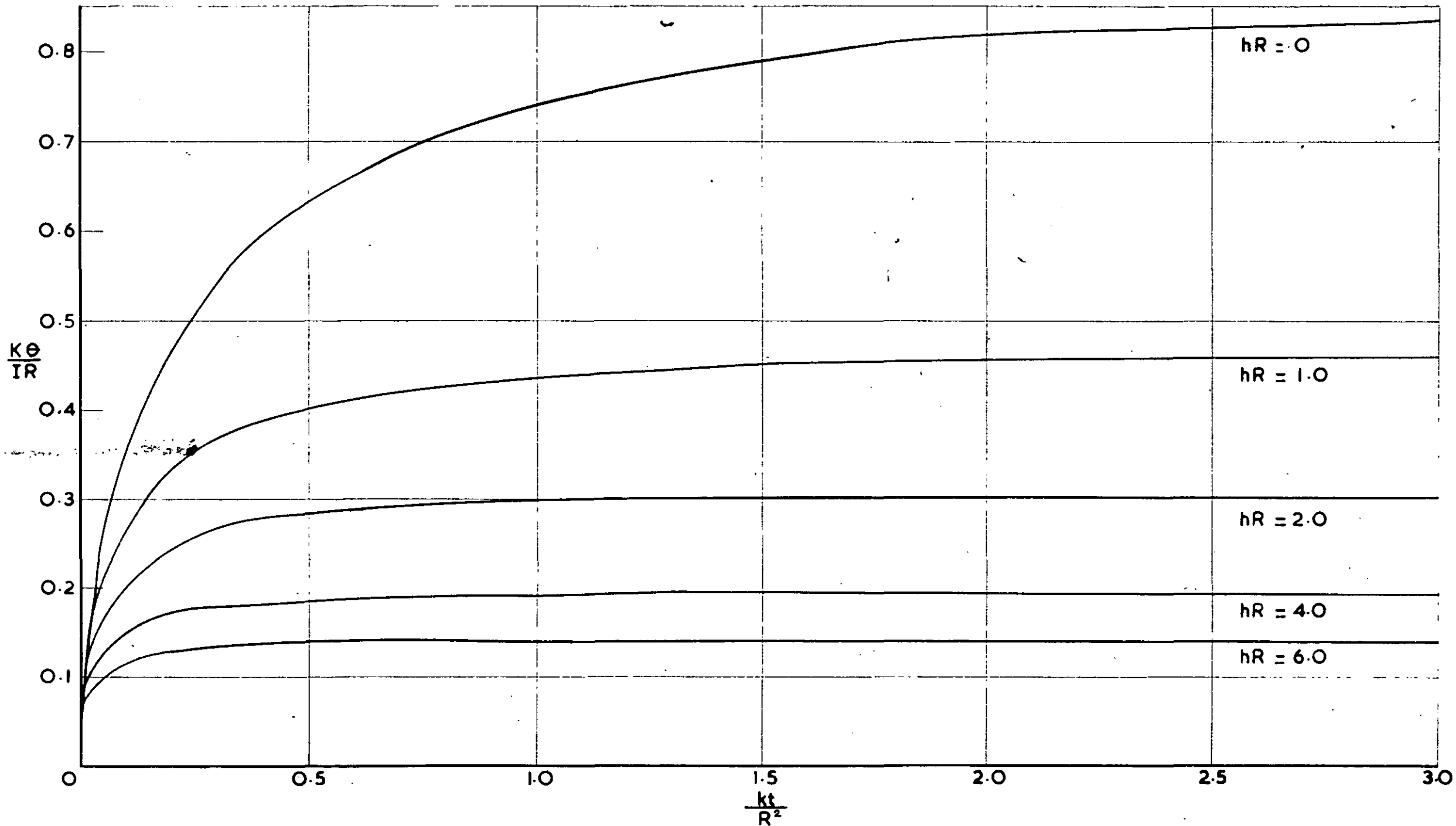


FIG. 1. TRANSIENT TEMPERATURE AT CENTRE OF HEATED DISK

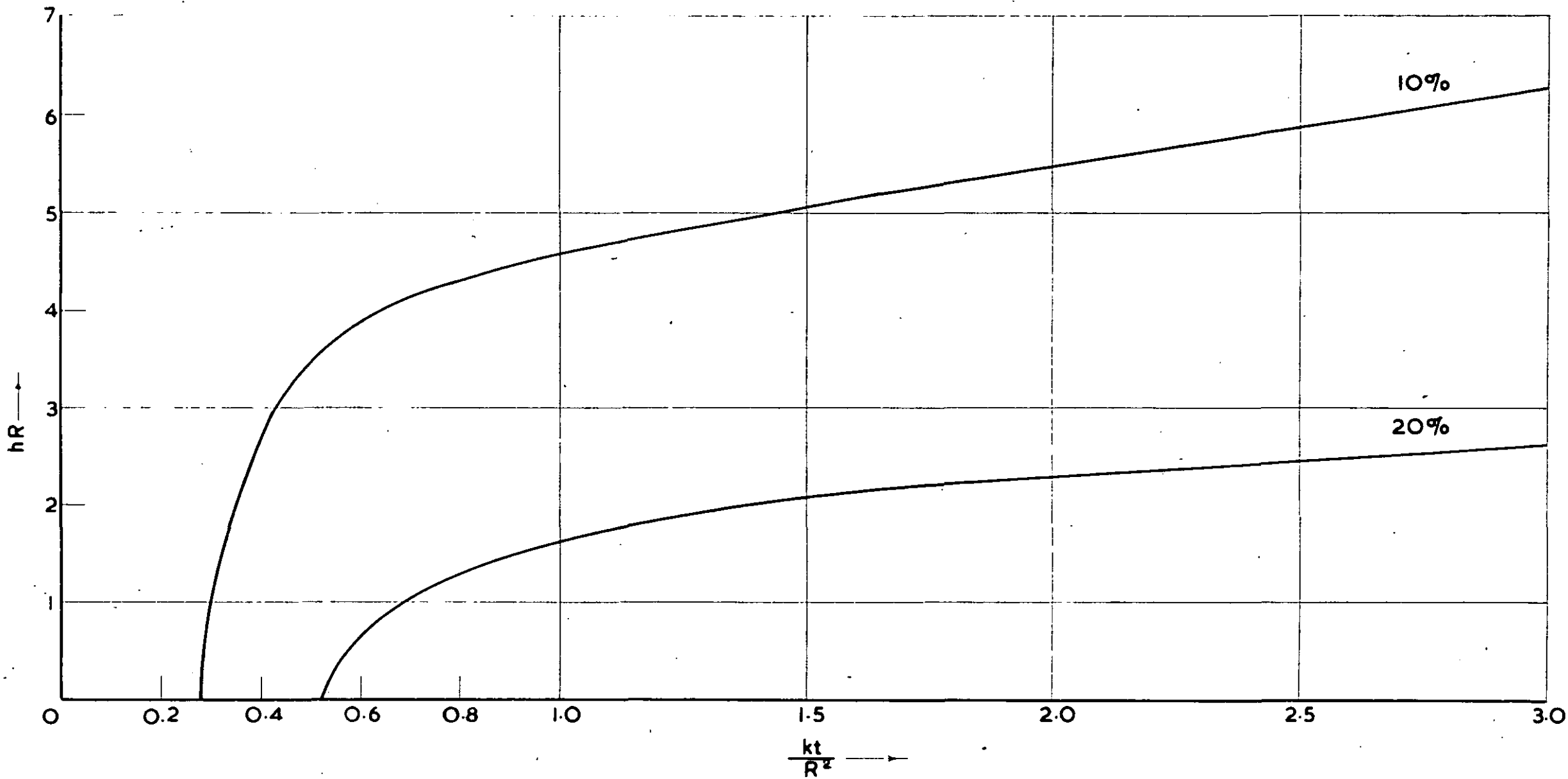


FIG. 2. VALUES OF  $\frac{kt}{R^2}$  AT WHICH TEMPERATURE OF DISK CENTRE IS 10% AND 20% BELOW THAT OF HEATED PLANE

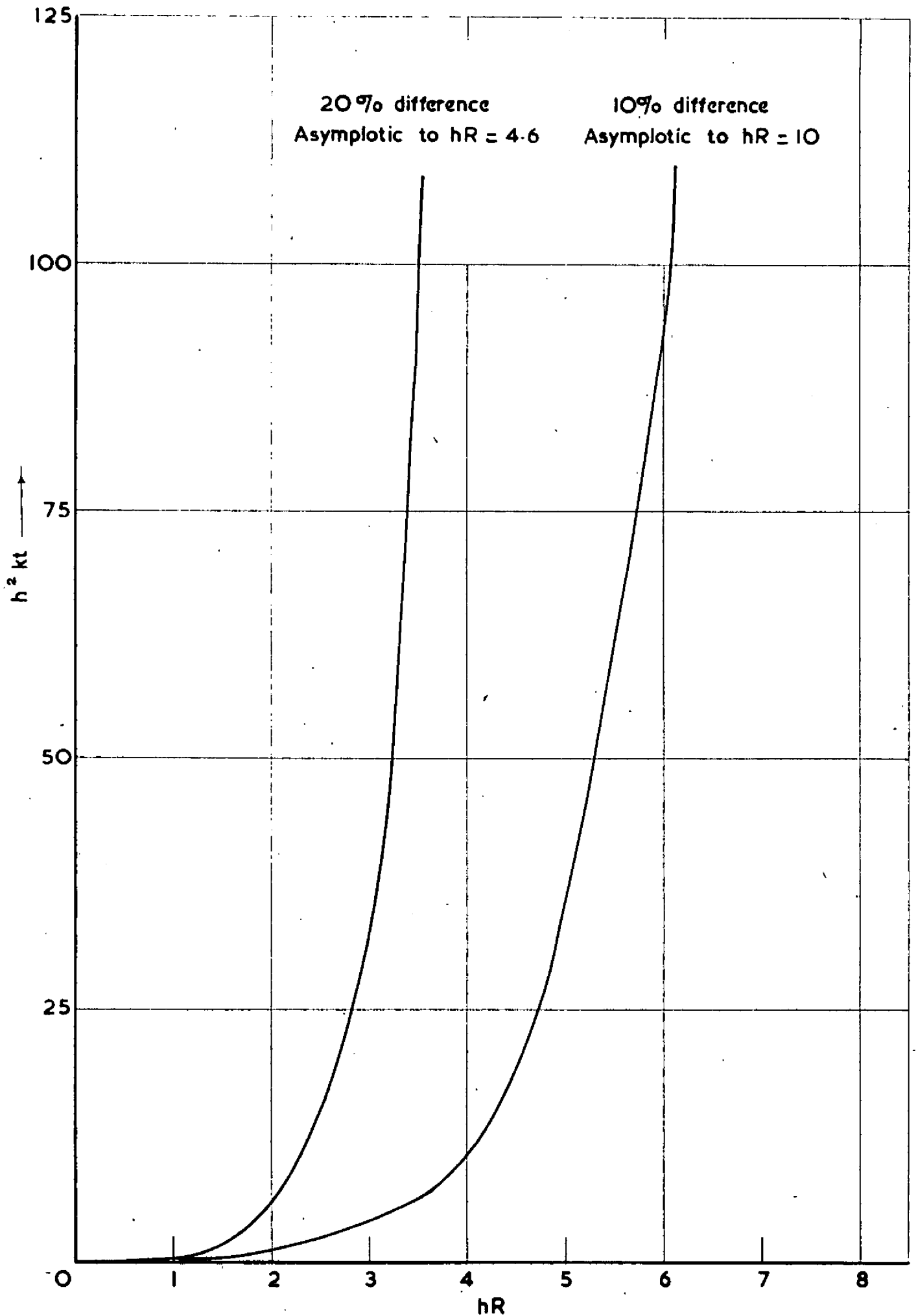


FIG. 3. VALUES OF  $h^2 kt$  AS FUNCTION OF  $hR$  FOR 10% & 20% DIFFERENCE BETWEEN DISK AND PLANE.